Positive temperature free fermions and solvable models in the **KPZ class**

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Plan of the talk

- Part 1. Presenting the problem we solve: connections between KPZ equation and determinantal/pfaffian point processes
- Part 2. Combinatorial construction of the correspondence KPZ models - free fermionic systems

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• Part 1. Presenting the problem we solve: connections between KPZ equation and determinantal/pfaffian point processes

KPZ equation



[Kardar-Parisi-Zhang '86]

 $\partial_t h = \frac{1}{2} \partial_x^2 h + \frac{1}{2} (\partial_x h)^2 + \eta$

• h = random height function $\eta =$ space-time white noise



KPZ equation [Kardar-F



[Kardar-Parisi-Zhang '86]

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- h = random height function $\eta =$ space-time white noise
- Well posedness: [Bertini-Giacomin '97], [Hairer '11], [Gubinelli-Perkowski-Imkeller '12]





$$\begin{cases} \partial_t h = \frac{1}{2}\partial_x^2 h + \frac{1}{2}(\partial_x h)^2 + \eta \\ h(x,0) = \log(\delta_x) \end{cases}$$



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Narrow wedge initial conditions

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[Amir-Corwin-Quastel, Calabrese-Le Doussal, Dotsenko, Sasamoto-Spohn '11]

$$\mathbb{E}\left[\exp\left(-ze^{h(0,t)+t/24}\right)\right] = \det\left(1 - fK_{\text{Airy}}\right)_{\mathscr{L}^{2}(\mathbb{R})}$$

$$f_{ry}(x,y) = \int_0^\infty \operatorname{Ai}(x+z)\operatorname{Ai}(y+z)dz \qquad f(x) = \frac{1}{1+e^{-xt^1}}$$

Airy Kernel

Fermi factor







$$\left| \begin{array}{l} \partial_t h = \frac{1}{2} \partial_x^2 h + \frac{1}{2} (\partial_x h)^2 + \eta & x \in \mathbb{R}_+ \\ h(x,0) = \log(\delta_x), & \partial_x h(x,t) \right|_{x=0} = A \end{array}$$



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 Narrow wedge initial conditions in half space

$$\left[\begin{array}{ll} \partial_t h = \frac{1}{2} \partial_x^2 h + \frac{1}{2} (\partial_x h)^2 + \eta & x \in \mathbb{R}_+ \\ h(x,0) = \log(\delta_x), & \partial_x h(x,t) \Big|_{x=0} = A \end{array} \right]$$

[Gueudre-Le Doussal'12, Borodin-Bufetov-Corwin'16, Barraquand-Borodin-Corwin-Wheeler'17 (*A=-1/2*), Krejenbrink-Le Doussal'19, De Nardis-Krejenbrink-Le Doussal-Thiery'20]

$$\mathbb{E}_{hs} \left[\exp\left(-ze^{h(0,t)+t/24}\right) \right] = \Pr\left(1-fK\right)_{\mathscr{L}^{2}(\mathbb{R})}$$

$$K(x,y) \qquad \qquad f(x) = \frac{1}{1+e^{-xt^{1/3}/z}}$$
2x2 matrix kernel
(Airy-like) Fermi factor

- Narrow wedge initial conditions
- Narrow wedge initial conditions in half space

- Mysterious relations between free fermions and KPZ equation
- Apparent only from the solutions
- Can we establish connections between KPZ eq. and free fermions a priori?

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- Solutions are obtained through Bethe Ansatz (BA)
- BA is very powerful but requires difficult calculations and (often) non-rigorous arguments
- Can we create an elementary theory to solve the KPZ eq.?









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$$h^{(q)}(x,t)$$

...how to solve discrete models?

$$\longrightarrow h^{KPZ}(x, t)$$

Typical model (FULL SPACE): q-Push TASEP [Borodin-Petrov '12]



 $h^{(q)}(x,t) = \eta_x(t)$



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 $\eta_n(t)\Big|_{t=0} = n$ • Typical feature: Assume step initial conditions



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Typical feature: Assume step initial control

$$\mathbb{P}(\eta_n(t) - n = k) = \sum_{\mu_1 = k} \frac{b_\mu \mathscr{P}_\mu(a)}{Z_{a_1}^q}$$

$$\bigoplus_{\substack{\eta_{n+2}}} \mathbb{Z} \qquad h^{(q)}(x,t) = \eta_x(t)$$

conditions
$$\eta_n(t)\Big|_{t=0} = n$$
 $+++\bullet\bullet\bullet\bullet$

 $,b_t$

q-Whittaker measure [Borodin-Corwin '11]

$$\mu = (\mu_1 \ge \mu_2 \ge \cdots \ge 0)$$





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$$\mathbb{P}(\eta_t^{\text{hs}}(t) - t = k) = \sum_{\mu_1 = k} \frac{b_{\mu}^{\text{el}} \mathscr{P}_{\mu}(t)}{Z_a^q}$$

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$$h^{(q)}(x,t) = \eta_x(t)$$



a;A

Half space *q*-Whittaker measure [BBC '18] $\mu = (\mu_1 \ge \mu_2 \ge \dots \ge 0)$





$$\mathbb{P}(\eta_n(t) - n = k) = \sum_{\mu_1 = k} \frac{b_\mu \mathscr{P}_\mu(a) \mathscr{P}_\mu(b_t)}{Z_{a,b_t}^q}$$

•KPZ behavior captured by marginal μ_1 of the q-Whittaker measures

 $\mathbb{P}(\eta_t^{\text{hs}}(t) - t = k) = \sum_{\mu_1 = k} \frac{b_{\mu}^{\text{el}} \mathscr{P}_{\mu}(a; A)}{Z_a^q}$



$$\mathbb{P}(\eta_n(t) - n = k) = \sum_{\mu_1 = k} \frac{b_\mu \mathscr{P}_\mu(a) \mathscr{P}_\mu(b_t)}{Z_{a,b_t}^q}$$

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techniques (Macdonald operators, Markov duality, Bethe Ansatz, Hecke



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- •KPZ behavior captured by marginal μ_1 of the q-Whittaker measures
- In *full space* case formulas for η_n or μ_1 are obtained through various algebras, Fock space...)
- **GOAL** : relate μ_1 with natural statistics of a determinantal/pfaffian point process

$$\mathbb{P}(\eta_t^{\text{hs}}(t) - t = k) = \sum_{\mu_1 = k} \frac{b_{\mu}^{\text{el}} \mathscr{P}_{\mu}(a)}{Z_a^q}$$

techniques (Macdonald operators, Markov duality, Bethe Ansatz, Hecke

In half space such techniques don't work (yet). Do pfaffian formulas exist?



Positive temperature Free Fermions

• Periodic Schur measure [Borodin '06]

$$\mathbb{P}(\lambda) = \frac{1}{\tilde{Z}_{a,b}^q} \sum_{\rho} q^{|\rho|} s_{\lambda/\rho} (d)$$

$$s_{\lambda/\rho}(a) = \det \left[h_{\lambda_i - \rho_j - i + j}(a) \right]_{i,j}$$
 Schur

- $\mathbb{P}(S = k) \propto q^{k^2/2} t^k$ for $k \in \mathbb{Z}$ independent of λ
- $(\lambda_1 + S, \lambda_2 + S, \lambda_3 + S, ...)$ is a determinantal point process

 $a)s_{\lambda/\rho}(b)$









Figures from [Betea-Bouttier]



Positive temperature Free Fermions



$$K(x, y) = \frac{1}{(2\pi i)^2} \oint_{|z|=r} \frac{dz}{z^{x+1}} \oint_{|w|=r'} \frac{dw}{w^{-y+1}} \frac{F(z)}{F(w)} \kappa(z, w),$$

$$F(z) = \prod_{i \ge 1} \frac{(b_i/z;q)_{\infty}}{(a_i z;q)_{\infty}} \qquad \qquad \kappa(z,w) = \sqrt{\frac{w}{z}} \frac{(q;q)_{\infty}^2}{(z/w,qw/z;q)_{\infty}} \frac{\vartheta_3(\zeta z/w;q)}{\vartheta_3(\zeta;q)}$$

$\mathbb{P}(\lambda_1 + S < s) = \det(1 - K)_{\ell^2\{s, s+1, \dots\}}$

• $(\lambda_1 + S, \lambda_2 + S, \lambda_3 + S, ...)$ is a determinantal point process with correlation





Free boundary Schur measure

• Free boundary Schur measure [Betea-Bouttier-Nejjar-Vuletic '17]

$$\mathbb{P}(\lambda) = \frac{\mathbf{1}_{\lambda' \text{even}}}{\tilde{Z}_{a;A}^{q}} \sum_{\rho' \text{even}} q^{|\rho|/2} s_{\lambda/\rho}($$

- $\mathbb{P}(S = k) \propto q^{2k^2} t^{2k}$ for $k \in \mathbb{Z}$ independent of λ
- $(\lambda_1 + 2S, \lambda_2 + 2S, \lambda_3 + 2S, ...)$ is a determinantal point process

(a;A)



Figure from [Betea-Bouttier]



 $L(x, y) = \begin{pmatrix} k(x, y) & -\nabla_y k(x, y) \\ -\nabla_x k(x, y) & \nabla_x \nabla_y k(x, y) \end{pmatrix}$

 $F(z) = \frac{(A/z;q)_{\infty}}{(Az;q)_{\infty}} \prod_{i>1} \frac{(a_i/z;q)_{\infty}}{(a_i z;q)_{\infty}}$

$\mathbb{P}(\lambda_1 + S < s) = Pf(1 - L)_{\ell^2\{s, s+1, ...\}}$

$k(x, y) = \frac{1}{(2\pi i)^2} \oint_{|z|=r} \frac{dz}{z^{x+3/2}} \oint_{|w|=r} \frac{dw}{w^{y+5/2}} F(z)F(w)\kappa^{hs}(z, w)$

$$\kappa^{\rm hs}(z,w) = \frac{(q,q,w/z,qz/w;q)_{\infty}}{(1/z^2,1/w^2,1/zw,qwz;q)_{\infty}} \frac{\vartheta_3(\zeta^2 z^2 w^2;q^2)}{\vartheta_3(\zeta^2;q^2)}$$



KPZ solvable models

Full space $\mu \sim \frac{b_{\mu} \mathscr{P}_{\mu}(x) \mathscr{P}_{\mu}(y)}{Z_{xy}^{q}}$

Half space $\mu^{hs} \sim \frac{b_{\mu}^{el} \mathscr{P}_{\mu}(a;A)}{Z_{a'A}^{q}}$

Determinantal/Pfaffian point processes

 $\lambda \sim \frac{1}{\tilde{Z}_{x,y}^q} \sum_{\rho} q^{|\rho|} s_{\lambda/\rho}(x) s_{\lambda/\rho}(y)$

 $\lambda^{\text{hs}} \sim \frac{I_{\lambda'\text{even}}}{\tilde{Z}_{a;A}^{q}} \sum_{\rho'\text{even}} q^{|\rho|/2} s_{\lambda/\rho}(a;A)$
KPZ solvable models



Half space $\mu^{hs} \sim \frac{b_{\mu}^{el} \mathscr{P}_{\mu}(a;A)}{Z_{a;A}^{q}}$

$$\mu_1 + \chi \stackrel{\mathcal{D}}{=} \lambda_1$$

 χ independent of μ_1 and $\mathbb{P}(\chi = n) = q^n (q^{n+1}; q)_{\infty}$

Determinantal/Pfaffian point processes

$$\lambda \sim \frac{1}{\tilde{Z}_{x,y}^q} \sum_{\rho} q^{|\rho|} s_{\lambda/\rho}(x) s_{\lambda/\rho}(y)$$

$$\lambda^{\text{hs}} \sim \frac{\mathbf{1}_{\lambda'\text{even}}}{\tilde{Z}_{a;A}^{q}} \sum_{\rho'\text{even}} q^{|\rho|/2} s_{\lambda/\rho}(a;A)$$

THEOREM (Imamura-M.-Sasamoto '21) $\mu_1^{\text{hs}} + \chi \stackrel{\mathcal{D}}{=} \lambda_1^{\text{hs}}$ (\star)



THEOREM (Imar $\mu_1 + \chi \stackrel{\mathcal{D}}{=} \lambda_1$

 χ independent of μ_1 and $\mathbb{P}(\chi = n) = q^n(q^n)$

Comments on (\star) :

- Reveals the origin of determinantal formulas for KPZ solvable models at positive temperature Nice symmetrical relation between full and half space
- Suggests a new paradigm to solve models
- Reveals combinatorial properties between Schur polynomials and q-Whittaker polynomials • Earlier results relating KPZ models and free fermions:
 - [Dean-Le Doussal-Majumdar-Schehr'15]
 - [Borodin'16], [Borodin-Gorin'16], [Borodin-Ohlshanki'16], [Borodin-Corwin-Barraquand-Wheeler'17]

nura-M.-Sasamoto '21)

$$\mu_1^{\text{hs}} + \chi \stackrel{\mathcal{D}}{=} \lambda_1^{\text{hs}} \qquad (\star)$$
ⁿ⁺¹; q)_m



Plan of the talk

• Part 2. combinatorial construction of the correspondence KPZ models - free fermionic systems

RSK.

• We prove (\star) combinatorially, developing a (bijective!) q-extension of the

- RSK.
- Combinatorial formulas



• We prove (\star) combinatorially, developing a (bijective!) q-extension of the

sum over semi-standard tableaux

- RSK.
- **Combinatorial formulas** \bullet





 $s_{\lambda/\rho}(x_1, x_2) = x_1^3 x_2 + x_1^2 x_2^2 + x_1^2 x_2^2 + x_1 x_2^3$

• We prove (\star) combinatorially, developing a (bijective!) q-extension of the

sum over semi-standard tableaux



- RSK.
- Combinatorial formulas



• We prove (\star) combinatorially, developing a (bijective!) q-extension of the

sum over "vertically strict tableaux"

$$\sum_{\lambda,\rho} q^{|\rho|} s_{\lambda/\rho}(x) s_{\lambda/\rho}(y) = \frac{1}{(q;q)_{\infty}} \prod_{i,j} \frac{1}{(x_i y_j;q)_{\infty}} \prod_{i,j} \frac{1}{(x_i y_j;q)_{\infty}}$$
$$\sum_{\mu} b_{\mu} \mathcal{P}_{\mu}(x) \mathcal{P}_{\mu}(y) = \prod_{i,j} \frac{1}{(x_i y_j;q)_{\infty}}$$

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$$(z;q)_{\infty} = \prod_{\ell \ge 0} (1 - q^{\ell} z)$$



 $(\star) \Leftrightarrow \sum_{\lambda,\rho} q^{|\rho|} s_{\lambda/\rho}(x) s_{\lambda/\rho}(y) = \sum_{\mu_1 + \nu_1 = k} q^{|\nu|} b_{\mu} \mathscr{P}_{\mu}(x) \mathscr{P}_{\mu}(y)$ $\lambda_1 = k$



$$\sum_{\lambda,\rho} q^{|\rho|} s_{\lambda/\rho}(x) s_{\lambda/\rho}(y) = \frac{1}{(q;q)_{\infty}} \prod_{i,j} \frac{1}{(x_i y_j;q_i)_{\infty}} \prod_{i,j} \frac{1}{(x_i y_j;q_j)_{\infty}} \sum_{\mu,\nu} q^{|\nu|} b_{\mu} \mathscr{P}_{\mu}(x) \mathscr{P}_{\mu}(y) = \frac{1}{(q;q)_{\infty}} \prod_{i,j} \frac{1}{(x_i;q_j)_{\infty}} \prod_{i,j} \frac{1}$$

$$(z;q)_{\infty} = \prod_{\ell \ge 0} (1 - q^{\ell} z)$$

 $\frac{1}{(q)_{\infty}} = \sum_{\lambda,\rho} \sum_{P,Q \in SST(\lambda/\rho)} q^{|\rho|} x^{P} y^{Q}$

 $(y_i; q)_{\infty}$

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$$\mathscr{K}(\mu) = \{\kappa = (\kappa_{1}, \dots, \kappa_{\mu_{1}}) : \kappa_{i} \ge \kappa_{i+1} \text{ if } \mu_{i}' = \mu_{i+1}'\}$$

$$b_{\mu} = \prod_{i \ge 1} (q;q)^{-1}_{\mu_{i}-\mu_{i+1}} = \sum_{\kappa \in \mathscr{K}(\mu)} q^{\kappa_{1}+\kappa_{2}+\dots+\kappa_{\mu_{1}}} \qquad \mu = \prod_{\kappa_{1} \ge \kappa_{2} \ge \kappa_{3}} \mu$$

$$\sum_{\lambda,\rho} q^{|\rho|} s_{\lambda l \rho}(x) s_{\lambda l \rho}(y) = \frac{1}{(q;q)_{\infty}} \prod_{i,j} \frac{1}{(x_{i} y_{j};q)_{\infty}} = \sum_{\lambda,\rho} \sum_{P,Q \in SST(\lambda l \rho)} q^{|\rho|} x^{P} y^{Q}$$

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$$= \{\kappa = (\kappa_{1}, ..., \kappa_{\mu_{1}}) : \kappa_{i} \ge \kappa_{i+1} \text{ if } \mu_{i}' = \mu_{i+1}'\}$$

$$b_{\mu} = \prod_{i\ge 1} (q;q)_{\mu_{i}-\mu_{i+1}}^{-1} = \sum_{\kappa\in\mathscr{K}(\mu)} q^{\kappa_{1}+\kappa_{2}+\dots+\kappa_{\mu_{1}}} \qquad \mu = \prod_{\kappa_{1}\ge\kappa_{2}\ge\kappa_{3}} \kappa_{4}\ge\kappa_{5}$$

$$(z;q)_{\infty} = \prod_{\ell \ge 0} (1 - q^{\ell} z)$$
$$(z;q)_{k} = \prod_{\ell=0}^{k-1} (1 - q^{\ell} z)$$

 $(\star) \Leftrightarrow$

$$\sum_{\substack{\lambda,\rho\\\lambda_1=k}} q^{|\rho|} s_{\lambda/\rho}(x) s_{\lambda/\rho}(y) = \sum_{\substack{\mu_1+\nu_1=k\\\mu_1+\nu_1=k}} q^{|\nu|} b_{\mu} \mathscr{P}_{\mu}(x) \mathscr{P}_{\mu}(y)$$





$$\sum_{\lambda,\rho} q^{|\rho|} s_{\lambda/\rho}(x) s_{\lambda/\rho}(y) = \frac{1}{(q;q)_{\infty}} \prod_{i,j} \frac{1}{(x_i y_j;q)_{\infty}} \prod_{i,j}$$

$$\sum_{\mu,\nu} \frac{q^{|\nu|} b_{\mu} \mathcal{P}_{\mu}(x) \mathcal{P}_{\mu}(y)}{(q;q)_{\infty}} \prod_{i,j} \frac{1}{(x_i y_j)_{\infty}} \frac{1}{(x_i y_j)_{\infty}} \sum_{i,j} \frac{1}{(x_i y_j)_{\infty}} \sum_{i$$

•
$$\mathscr{K}(\mu) = \{\kappa = (\kappa_1, \dots, \kappa_{\mu_1}) : \kappa_i \ge \kappa_{i+1} \text{ if } \mu_i'\}$$

$$(z;q)_{\infty} = \prod_{\ell \ge 0} (1 - q^{\ell} z)$$
$$(z;q)_{k} = \prod_{\ell=0}^{k-1} (1 - q^{\ell} z)$$

 $\frac{1}{(q)_{\infty}} = \sum_{\lambda,\rho} \sum_{P,Q \in SST(\lambda/\rho)} q^{|\rho|} x^{P} y^{Q}$ $\frac{1}{y_j;q)_{\infty}} = \sum_{\mu,\nu} \sum_{\kappa \in \mathcal{K}(\mu)} \sum_{V,W \in VST(\mu)} q^{|\nu|+|\kappa|+\mathcal{H}(V)+\mathcal{H}(W)} x^V y^W$ $i'_{i} = \mu'_{i+1}$

IDEA: $(P, Q) \xleftarrow{\Upsilon} (V, W; \kappa, \nu)$

 $(\star) \Leftrightarrow \sum_{\lambda} q^{|\rho|} s_{\lambda/\rho}(x) s_{\lambda/\rho}(y) = \sum_{\lambda} q^{|\nu|} b_{\mu} \mathscr{P}_{\mu}(x) \mathscr{P}_{\mu}(y)$ $\mu_1 + \nu_1 = k$ λ, ρ $\lambda_1 = k$





Construction of $(P, Q) \leftrightarrow (V, W; \kappa, \nu)$ (sketch)

Construction of $(P, Q) \leftrightarrow (V, W; \kappa, \nu)$ (sketch) $P \qquad Q$







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• First construct u by "squeezing" P, Q



Construction of $(P, Q) \leftrightarrow (V, W; \kappa, \nu)$ (sketch) $P \qquad Q$





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 $\nu =$



Construction of $(P, Q) \leftrightarrow (V, W; \kappa, \nu)$ (sketch) $\nu =$

• First construct ν by "squeezing" P, Q





- First construct ν by "squeezing" P, Q
- From now on assume pair (P, Q) is "squeezed"






































































































































































• To construct (V, W) we define a dynamics on tableaux based on internal insertion [Sagan-Stanley'90]

$(P, Q) \rightarrow (P', Q')$









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 $(P, Q) \rightarrow (P', Q') \rightarrow \cdots \rightarrow (P^{(n)}, Q^{(n)})$



RSK









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 $(P,Q) \rightarrow (P',Q') \rightarrow \cdots \rightarrow (P^{(n)},Q^{(n)}) \longrightarrow (V,W)$







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W_{-}				
	1	2	2	3
	2	5	Ŋ	
	3			

 \mathcal{T} (0,2,1,1)



• $(P, Q) \leftrightarrow (V, W; \tau)$ is a bijection, but τ depends on (V, W)!



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LEMMA (IMS'21) $\tau_i = \kappa_i + \mathcal{H}_i(V$ $\mathcal{H}_i = \text{local energy function}$



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Comments on (Δ) :

- integrable systems, Kashiwara crystals, ...
- \mathscr{H}_i appear in the description of *phase shift* in the \mathfrak{Sl}_n Box and Ball System



$$(V) + \mathcal{H}_i(W)$$

Shows deep connections between theories of skew tableaux and theories of discrete

• This is because the skew RSK dynamics is a generalization of the Box and Ball Systems





- Classical RSK commutes with \mathfrak{Sl}_n Kashiwara operators
- Energy \mathscr{H} comes from affine $(\widehat{\mathfrak{Sl}}_n)$ crystals

• To prove energy (\mathscr{H}) formulas we need to study symmetries of the skew RSK





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• To prove energy (\mathscr{H}) formulas we need to study symmetries of the skew RSK

• IDEA : equip (P, Q) and (V, W) of $\widehat{\mathfrak{gl}}_n$ - bicrystal structure preserved by RSK



Affine crystal graphs

- Vertically strict tableaux are tensor product of KR crystals
- The crystal graph $VST(\mu)$ is connected [Kashiwara'90]
- Demazure subgraph: remove "bad" $\widetilde{e}_0, \widetilde{f}_0$ arrows
- For any V there exist a path on the Demazure subgraph $\mathscr{L}_V : V \mapsto \mathrm{YAM}(\mu)$

•
$$\mathscr{H}(V) = \# \tilde{f}_0 - \# \tilde{e}_0$$
 in \mathscr{L}_V

[Schilling-Tingley'12]

 $\begin{array}{c|c}2 & 1 & 1\\\hline 3 & \end{array}$



Affine Bicrystal (P, Q)

- On (V, W) bicrystal structure is product structure
- Impose $(P, Q) \mapsto (V, W)$ preserve bicrystal structure
- On (V, W) bicrystal structure is NOT product structure (nontrivial \tilde{e}_0, \tilde{f}_0)
- We transport maps $\mathscr{L}_V, \mathscr{L}_W \mapsto \mathscr{L}_{P,Q}$

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LEMMA (IMS'21)



 $\mathscr{L}_{P,O}$ linearizes the skew RSK map



To sum up

- $(P, Q) \longleftrightarrow (V, W; \kappa, \nu)$ has the following properties
 - P and V have equal content. Same for Q and W
 - $|\rho| = \mathcal{H}(V) + \mathcal{H}(W) + |\kappa| + |\nu|$ $(\mathcal{H} = \mathcal{H}_1 + \mathcal{H}_2 + \cdots)$
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 $- |\nu|$ $(\mathcal{H} = \mathcal{H}_1 + \mathcal{H}_2 + \cdots)$


Conclusion

- Our bijective theory gives direct connection between solvable non-free fermionic models and positive temperature free fermionic models
- Mysterious appearence of a discrete integrable system (KdV like models)
- We construct a bijective q-extension of the RSK correspondence Υ
- With Υ we can prove bijectively the Cauchy identities (CI) for q-Whittaker polynomials (first time)
- With Υ we can prove bijectively refinements of the CI relating q-Whittaker and skew-Schur polynomials